

A detailed map of the Cosmic Microwave Background (CMB) showing temperature fluctuations across the sky. The map is a complex, noisy pattern of blue and orange/yellow colors, representing the distribution of matter and energy in the early universe. The colors are more concentrated in the upper half of the image, with the lower half being a solid, dark grey color.

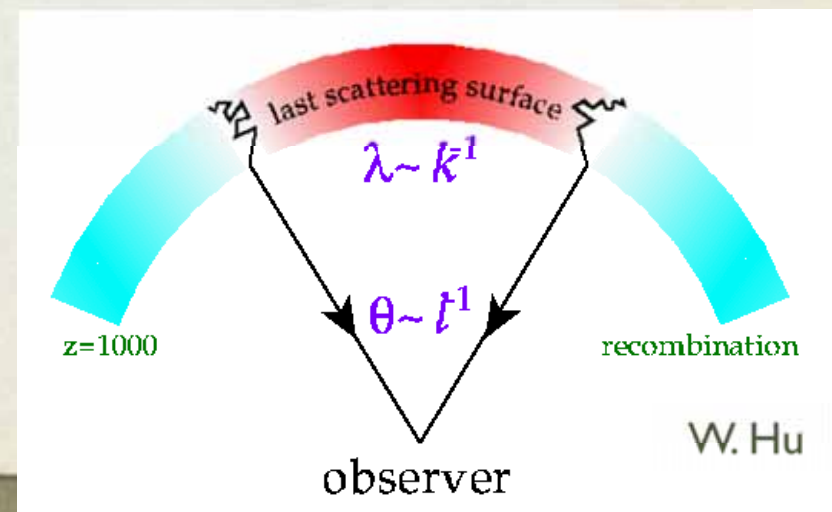
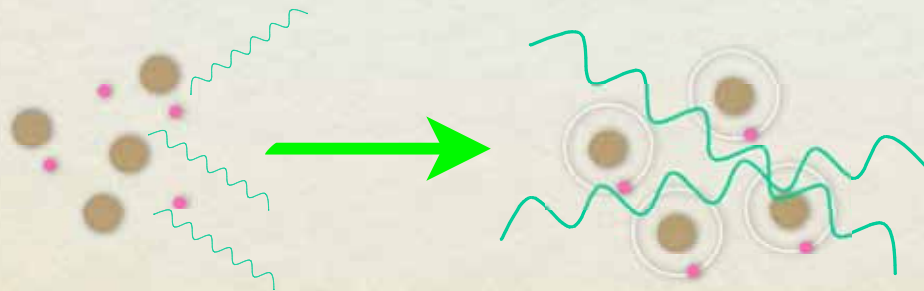
PHYSICS OF THE CMB

Jonathan Pritchard

THE COSMIC MICROWAVE BACKGROUND

- 400,000 years after the Big Bang, the temperature of the Universe was $T \sim 10,000$ K
- Hot enough to keep hydrogen atoms *ionized* until this time
 - *proton + electron* \rightarrow *Hydrogen* + *photon* [$p^+ + e^- \rightarrow H + \gamma$]
 - *charged plasma* \rightarrow *neutral gas*
- Photons (light) can't travel far in the presence of charged particles

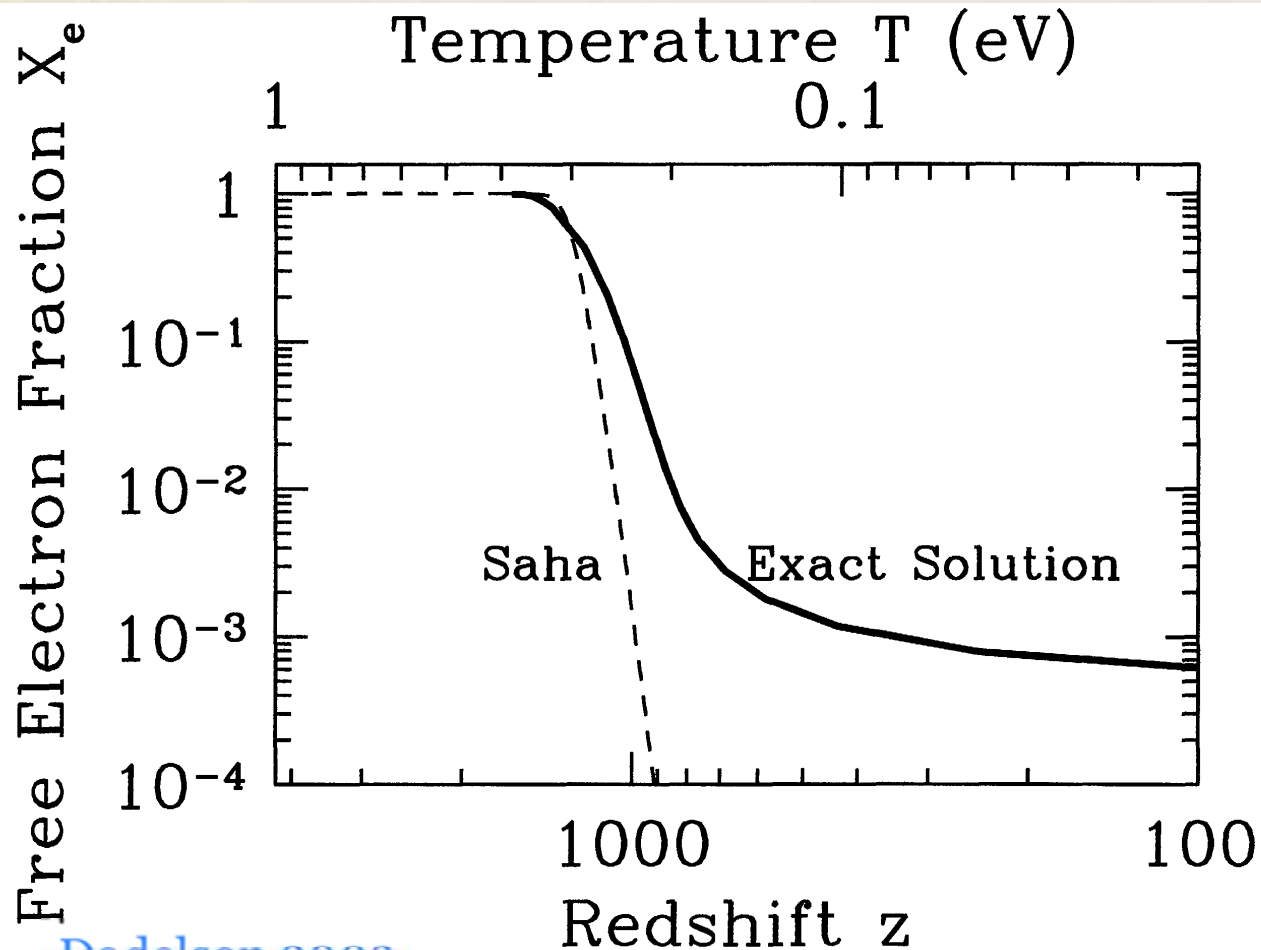
• *Opaque* \rightarrow *transparent*



RECOMBINATION

Saha Equation

$$\frac{1 - X_e}{X_e^2} = \frac{4\sqrt{2}\zeta(3)}{\sqrt{\pi}} \eta \left(\frac{kT}{m_e c^2} \right)^{3/2} e^{B/(kT)}$$



Recombination
once $T \sim 0.3$ eV



Meghnad Saha

DECOUPLING

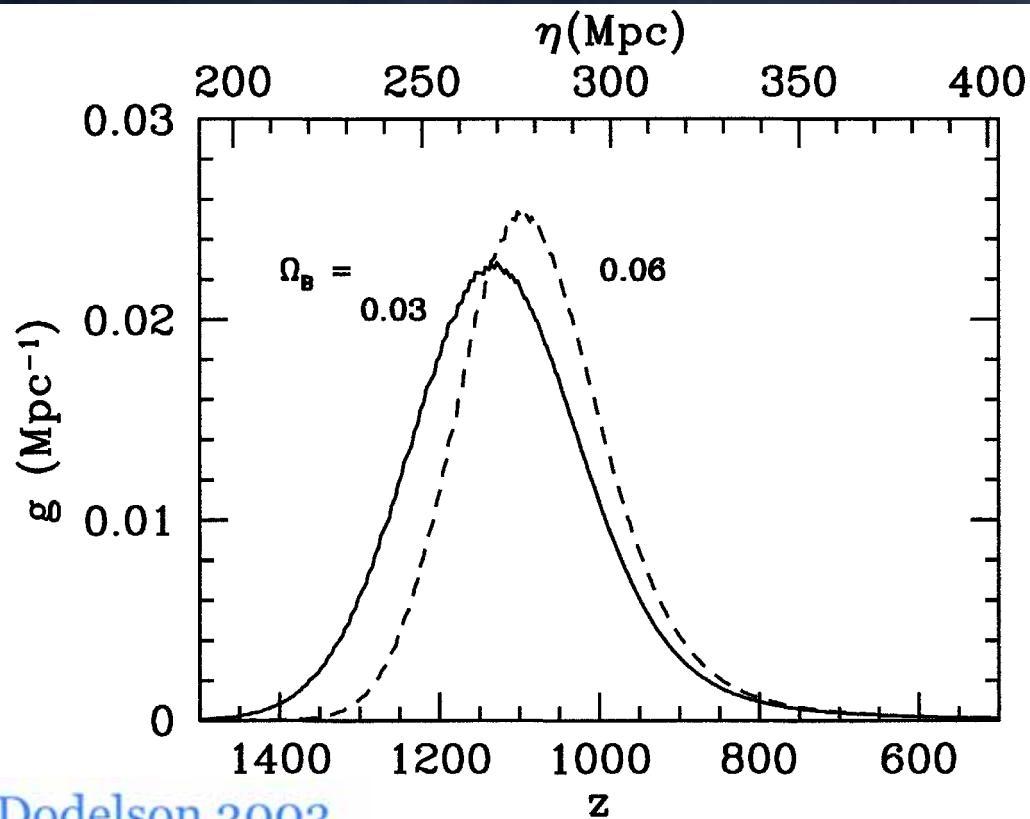
$$\tau = \int \sigma_T n_e(z) dl$$

Optical depth

$$g(\eta) \equiv -\dot{\tau} e^{-\tau}$$

visibility function

$$\int_0^{\eta_0} d\eta g(\eta) = 1,$$

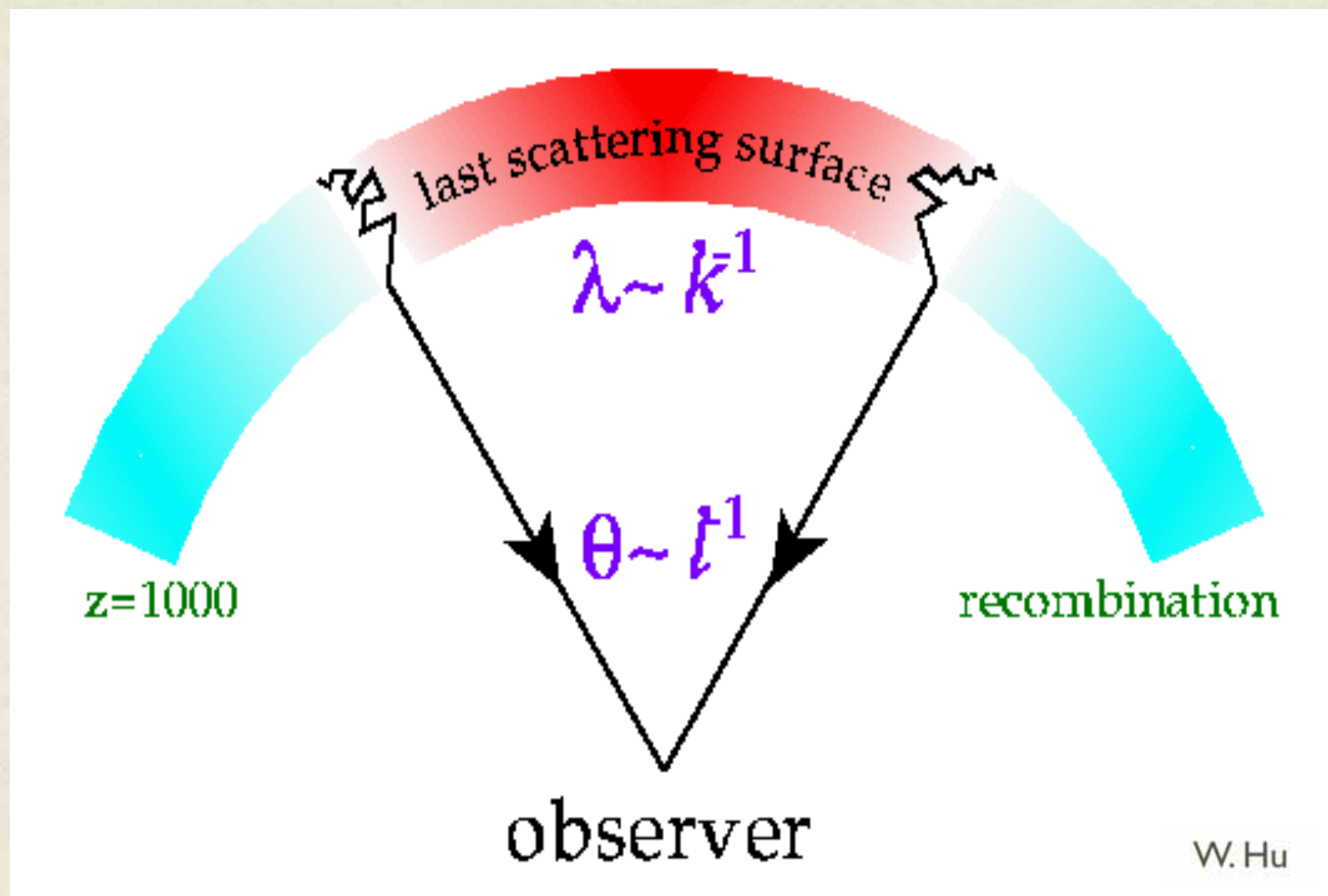


Dodelson 2003

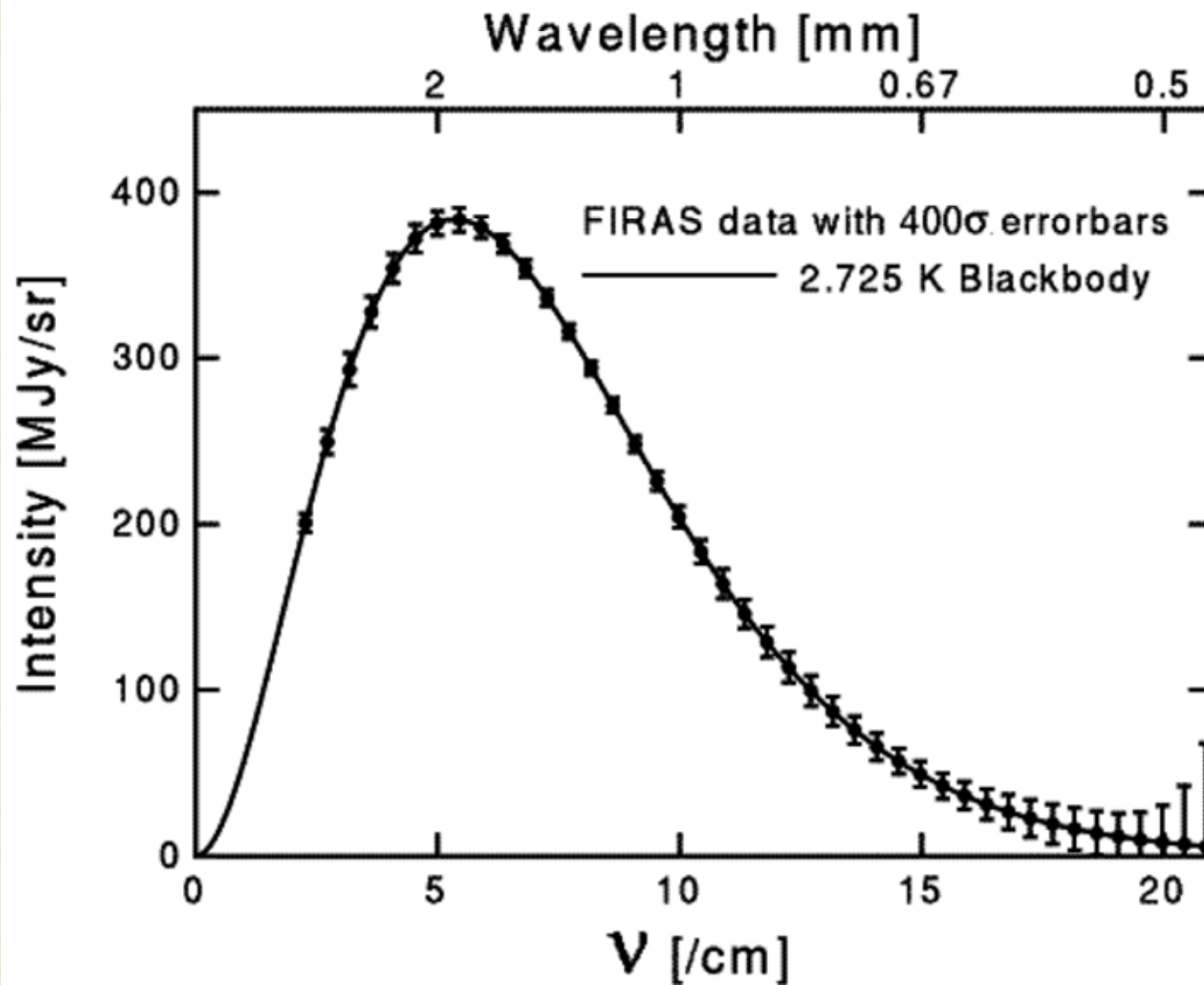
Visibility defines the
duration of decoupling
~PDF for last scattering

Peak at $z \sim 1100$

SURFACE OF LAST SCATTERING



CMB BLACKBODY



COBE-FIRAS
(1994)

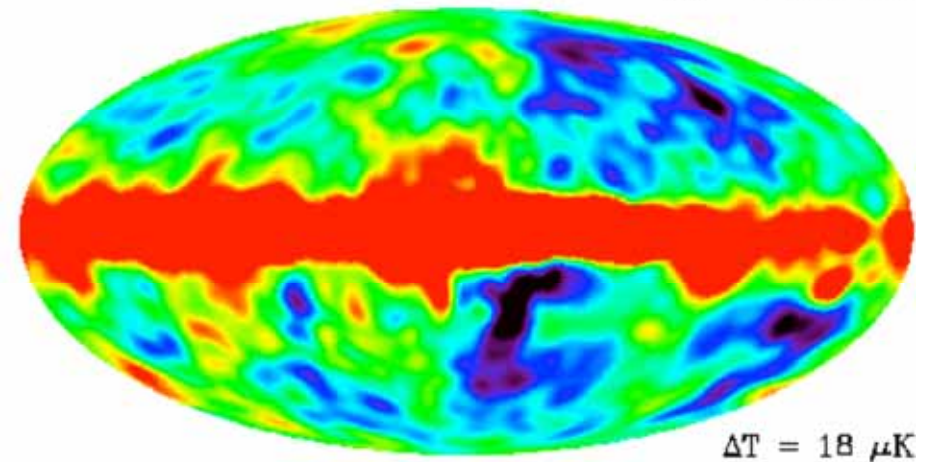
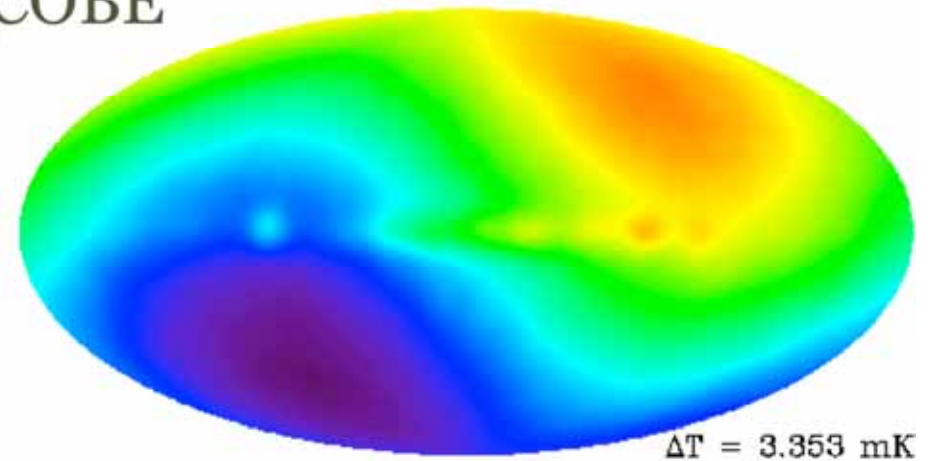
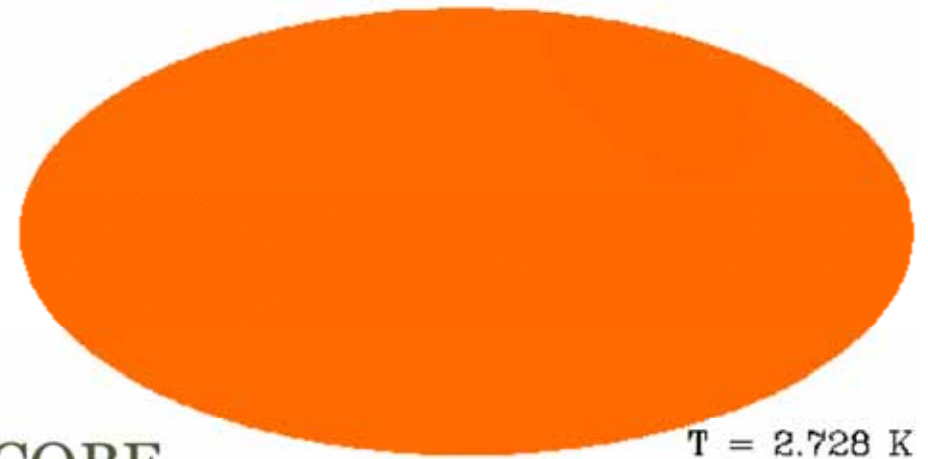
Strong evidence
that Universe was
thermalised at
decoupling

CMB TEMPERATURE

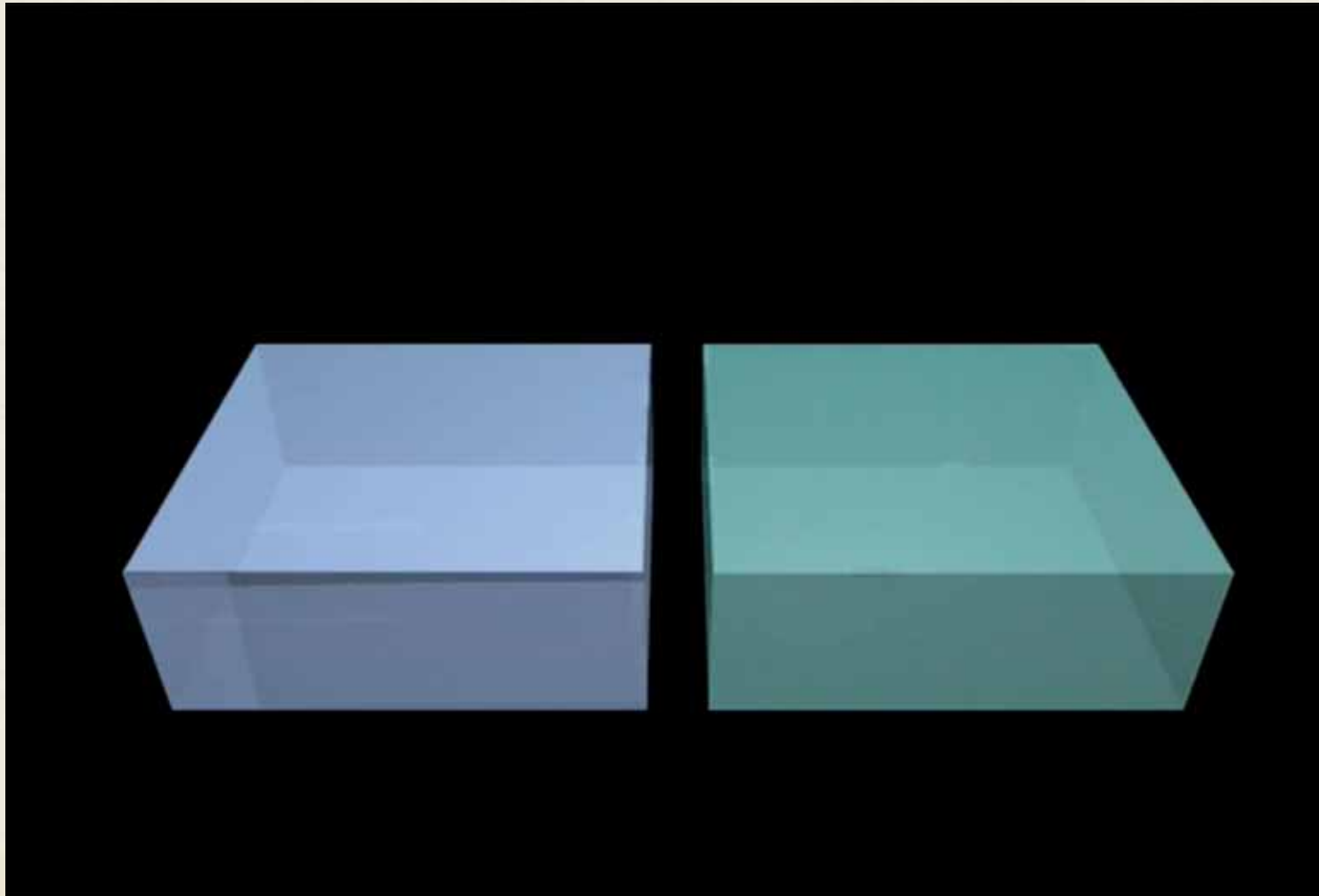
CMB temperature close to uniform

Motion relative to CMB rest frame
leads to dipole pattern

Temperature fluctuations at
level of $\sim 10 \mu\text{K}$ seen



TEMPERATURE FLUCTUATIONS



MULTIPOLE DECOMPOSITION

Natural to decompose temperature fluctuations in angular basis

$$\frac{T(\hat{x}) - \bar{T}}{\bar{T}} \equiv \frac{\Delta T}{T}(\hat{x}) = \sum_{\ell m} a_{\ell m} Y_{\ell m}(\hat{x})$$

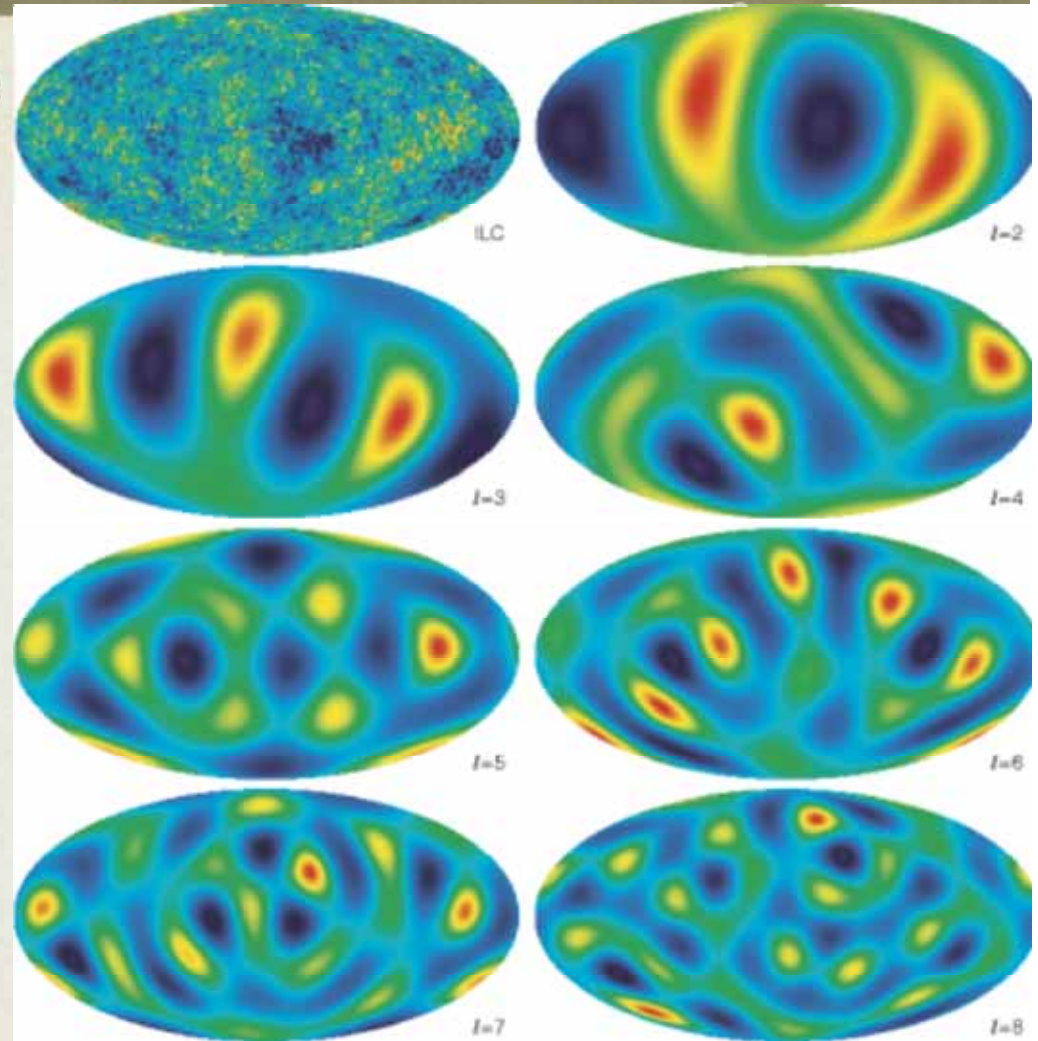
$Y_{\ell m}$ are spherical harmonics

Angular multipole $l \sim \pi/\theta$

Power spectrum:

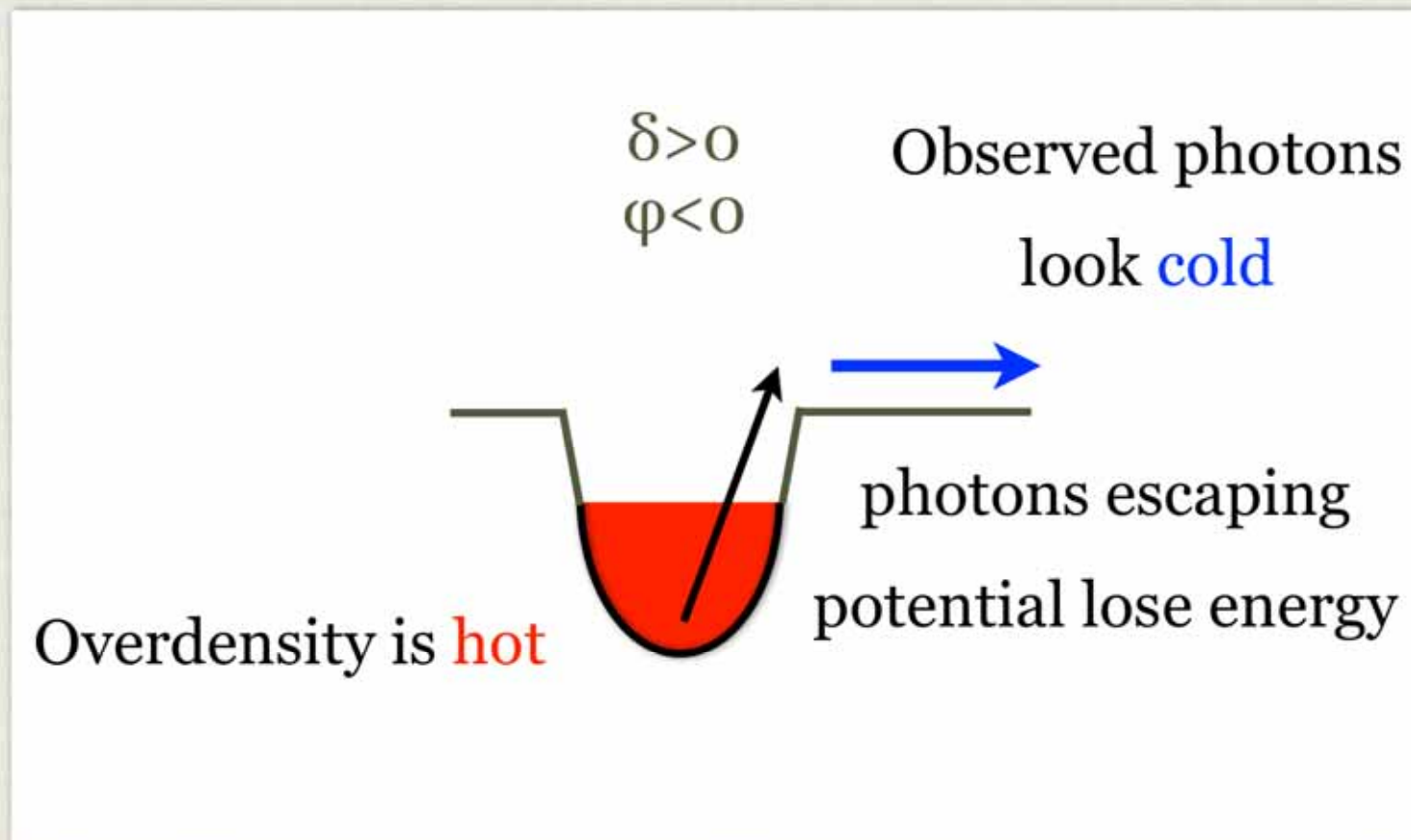
$$C_\ell = \langle |a_{\ell m}|^2 \rangle$$

For Gaussian fluctuations, power spectrum contains all information



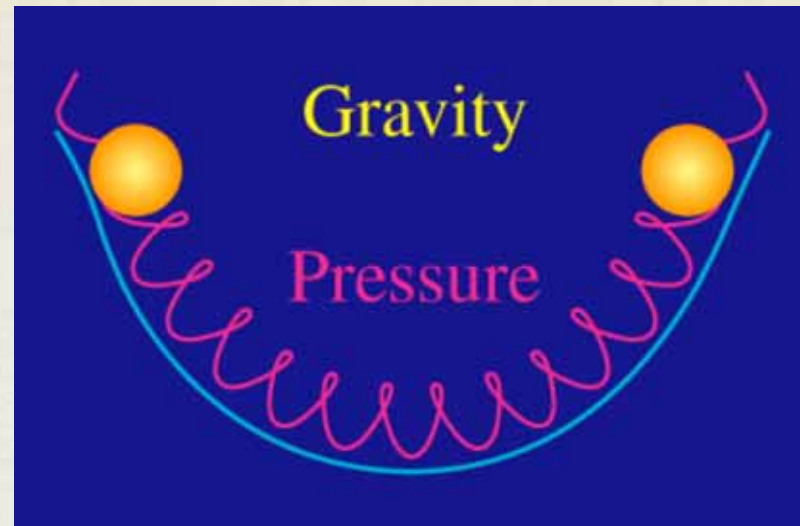
SACHS-WOLFE EFFECT

- Overdense regions will be intrinsically hotter BUT
See light after it escapes gravitational well



- observed **cold** spots = **overdense** regions at last scattering
observed **hot** spots = **underdense** regions at last scattering

PRESSURE VS GRAVITY



- Scale outside horizon - growing under gravity
- On scales inside horizon - pressure balances gravity
 - => oscillation
 - => regions denser or less dense

Basic picture is that of forced harmonic oscillation

$$\ddot{x} + \frac{k}{m}x = \frac{F_0}{m}.$$

- 1) Evolving potentials act as forcing term
- 2) Baryon content determines “mass”

$$c_s \equiv \sqrt{\frac{1}{3(1+R)}} \quad R = \frac{3\Omega_b h^2}{4\Omega_\gamma h^2}$$

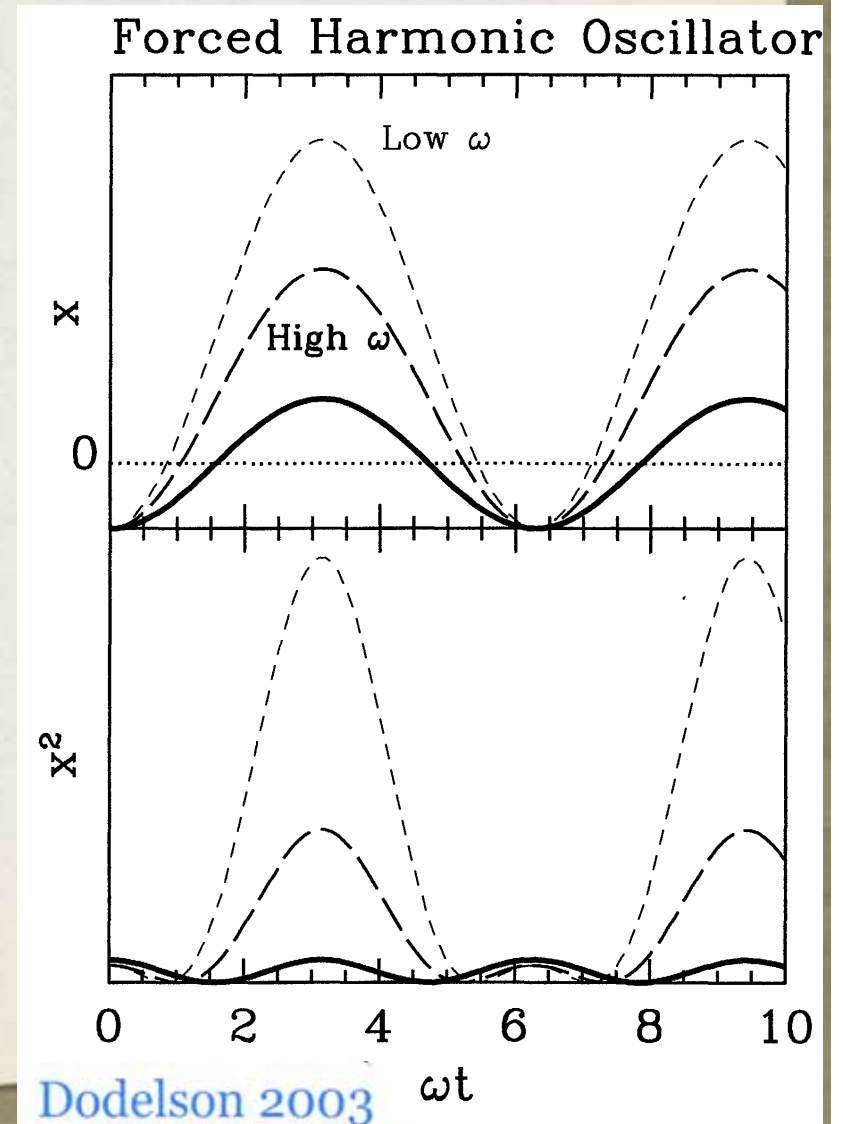
Solution:

$$x = A \cos(\omega t) + \frac{F_0}{m\omega^2}.$$

Power spectrum from a_{lm}^2

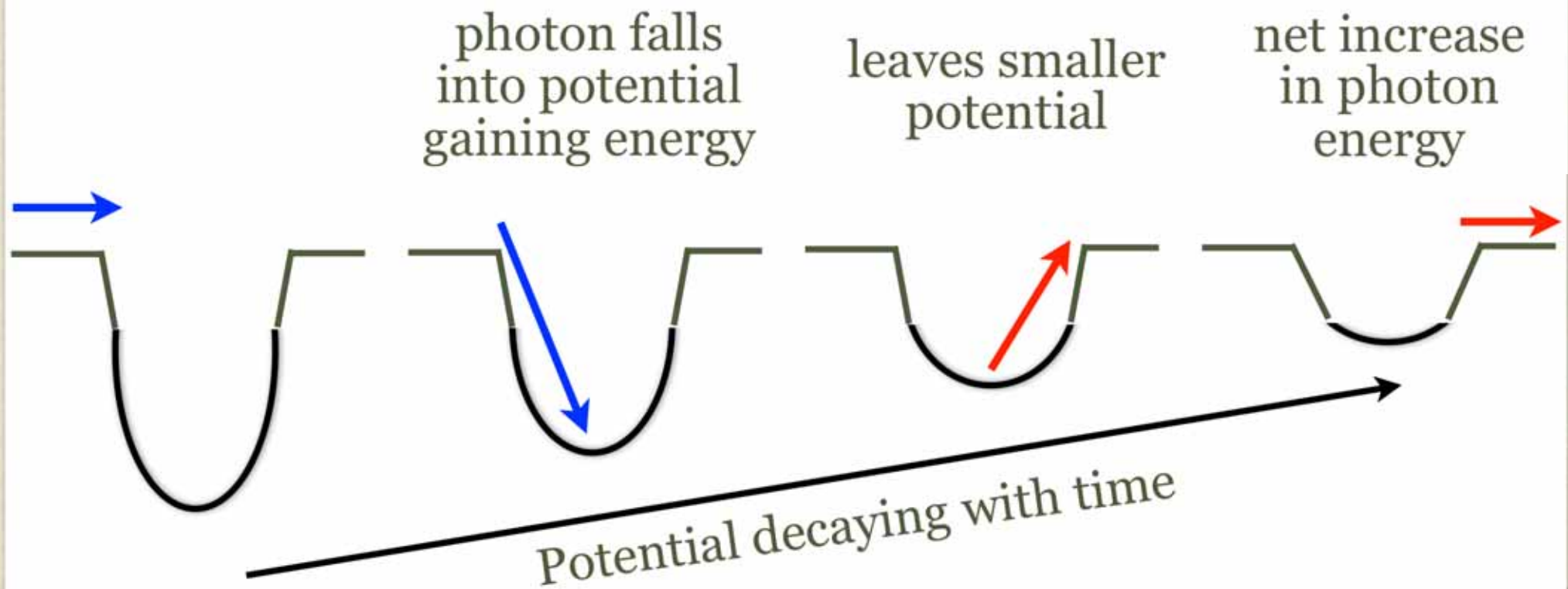
=> odd peaks (compression)

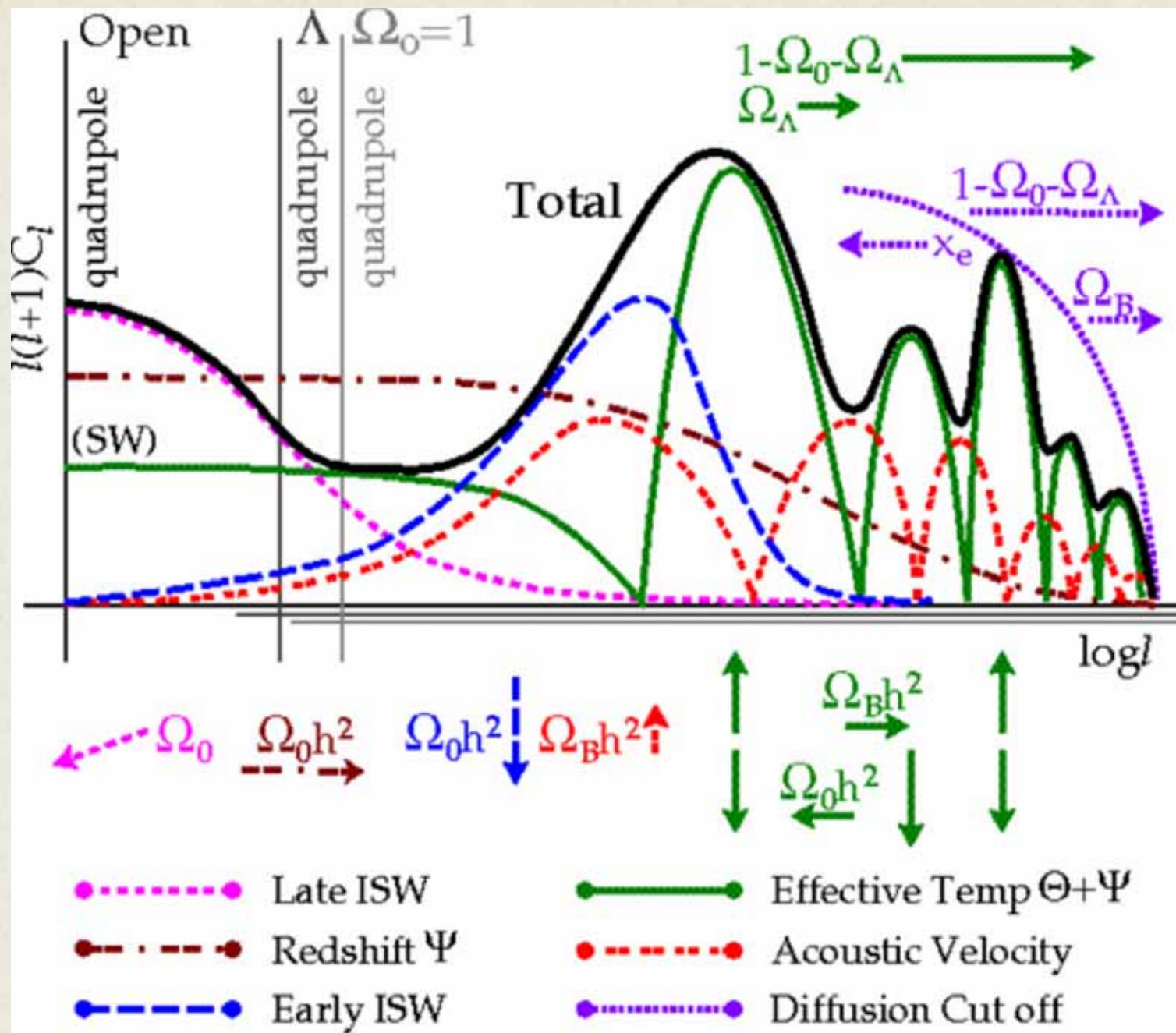
more enhanced than even (rarification)



INTEGRATED SACHS-WOLFE EFFECT

- A photon gains energy moving through a decaying potential
=> boost to height of first peak
- For later peaks oscillations average out gain in energy
- Early ISW as radiation domination ends and matter domination starts
Late ISW effect when dark energy becomes important



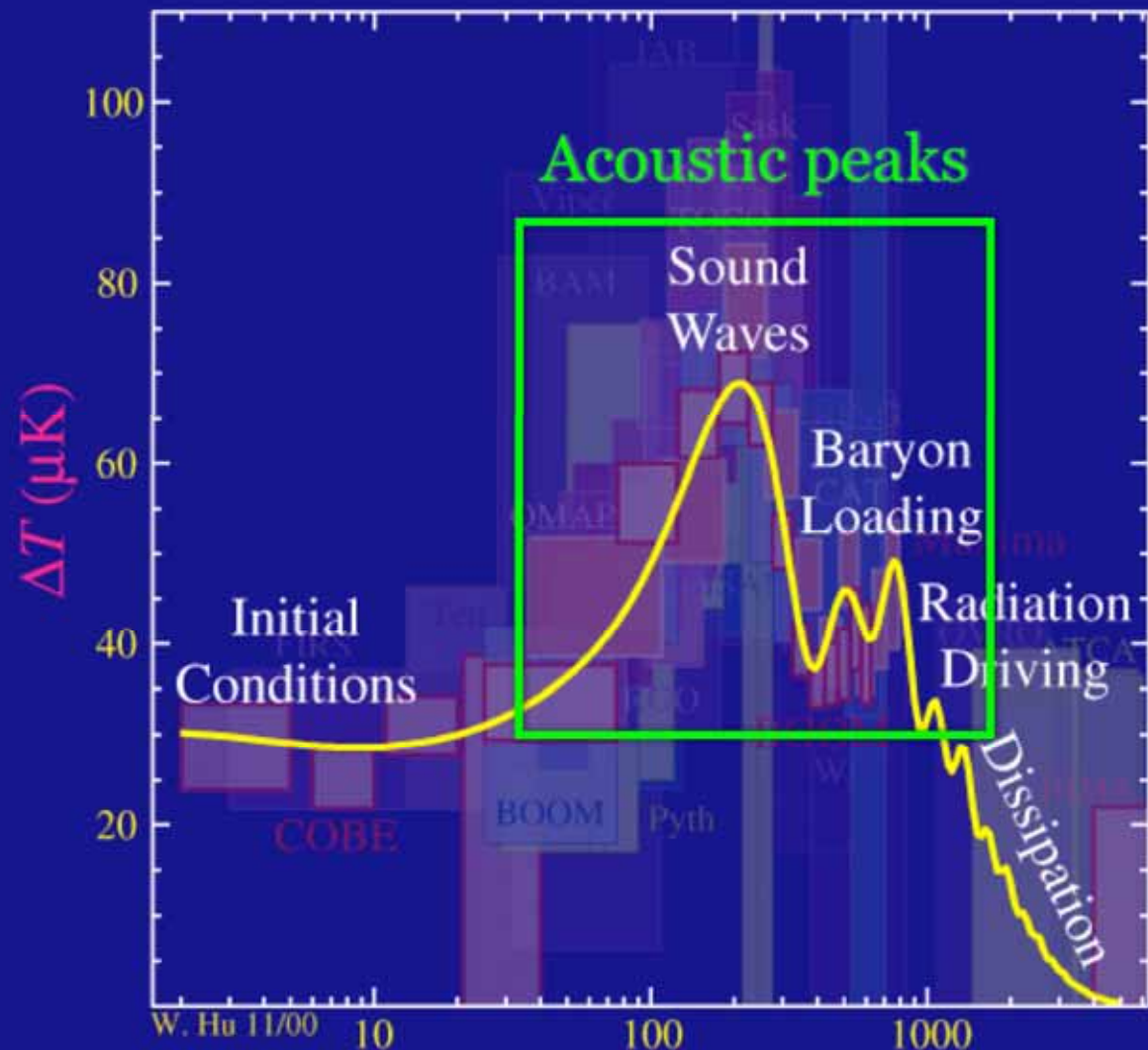


ANGULAR POWER SPECTRUM

Scales inside horizon
are oscillating
=> peaks are different
harmonics of oscillation

“Acoustic peaks”

More baryons
=> lower sound speed
=> higher third peak



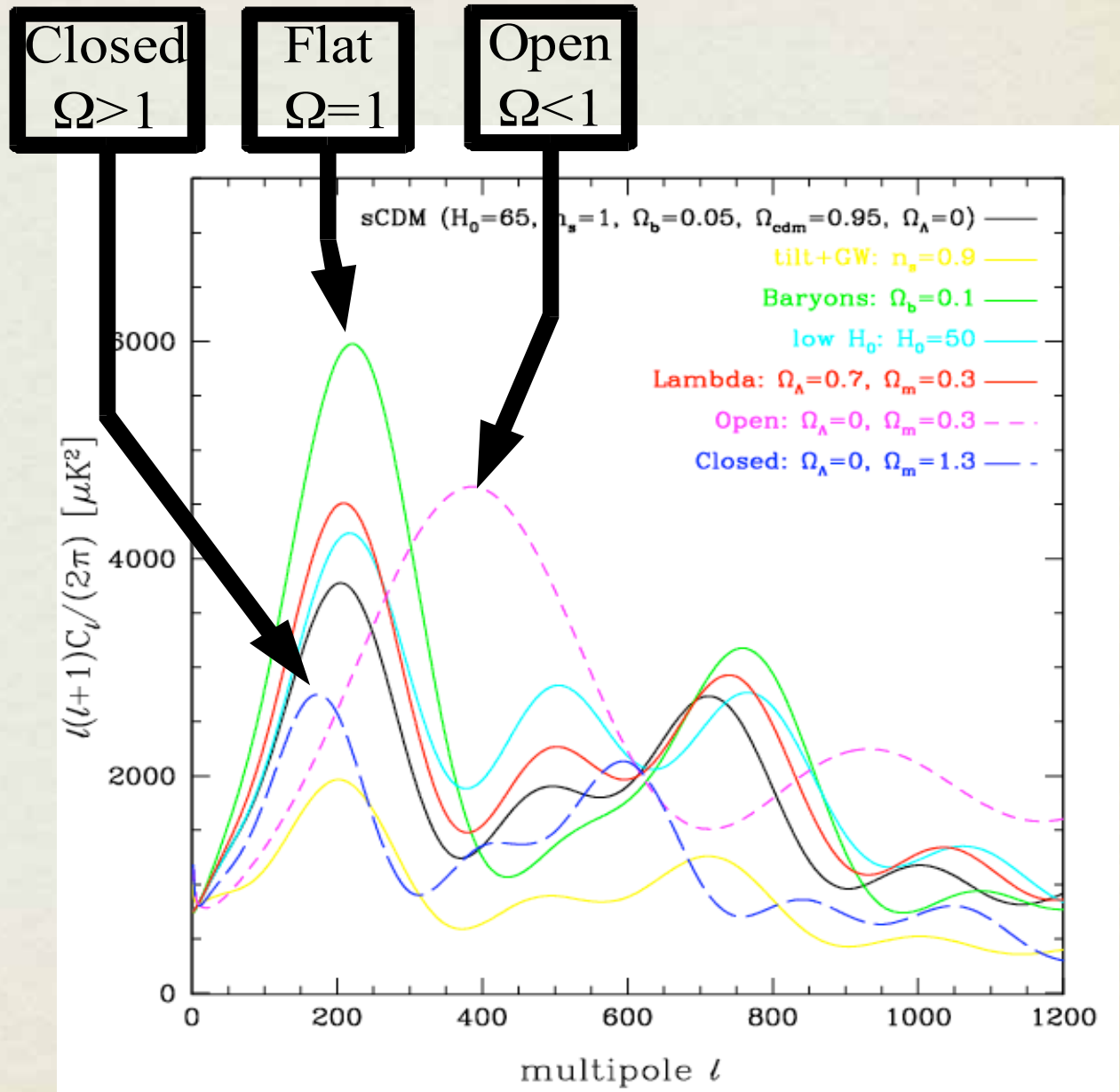
- $\ell_{\text{peak}} \sim \pi/\theta_{\text{fluc}}$
 - (Location of peak in C_ℓ)
 - $\sim (\pi/\text{Angular scale of typical fluctuations})$

- First peak at sound horizon/angular diameter distance

$$\theta_{\text{MC}} \approx r_s(z_{\text{CMB}})/D_A(z_{\text{CMB}})$$

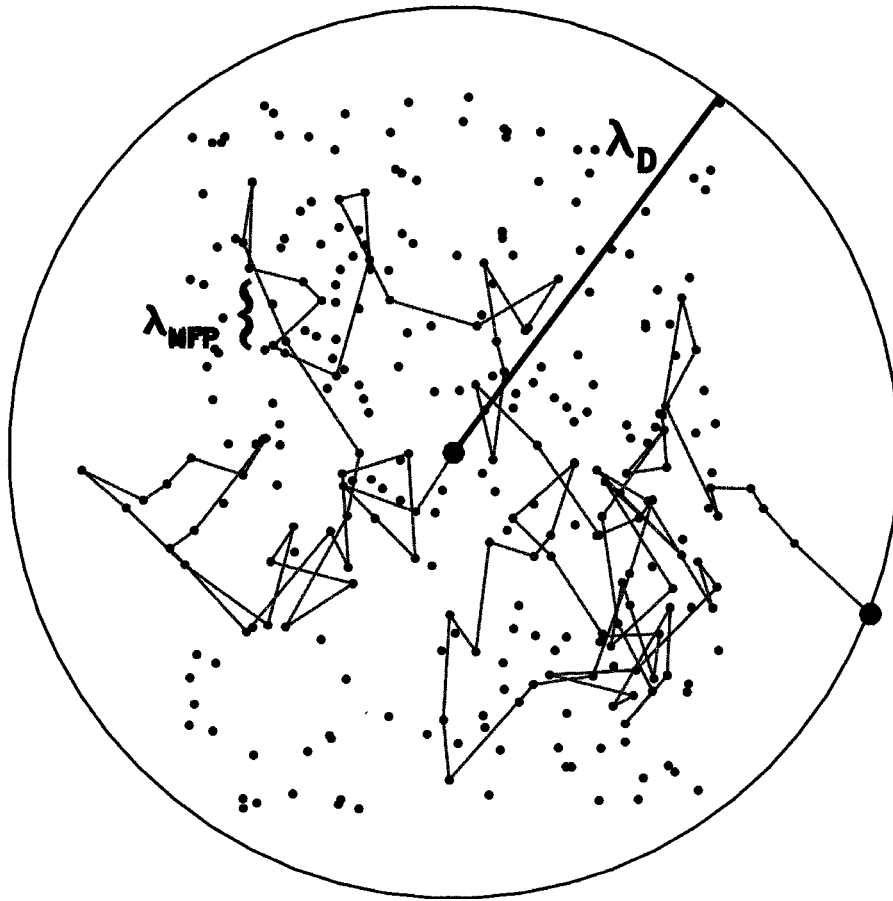
(Approximately)
sensitive to *curvature*

First peak and curvature



DIFFUSION AND DAMPING

Photon Diffusion



Recombination is not instantaneous => gives photons time to diffuse

$$\lambda_D \sim \lambda_{\text{MFP}} \sqrt{n_e \sigma_T H^{-1}}$$
$$= \frac{1}{\sqrt{n_e \sigma_T H}}$$

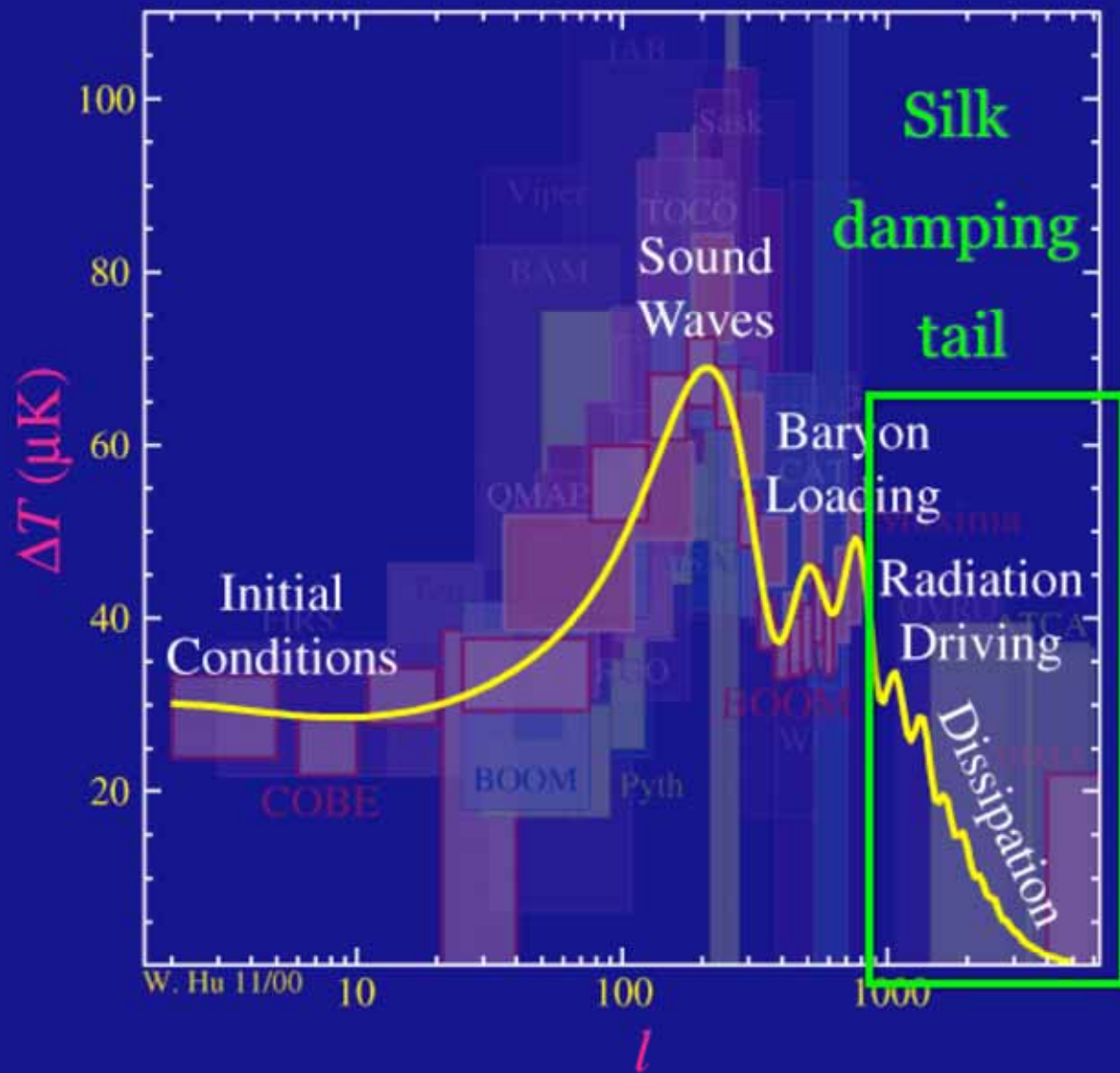
Photons move between hot and cold regions so fluctuations average out

ANGULAR POWER SPECTRUM

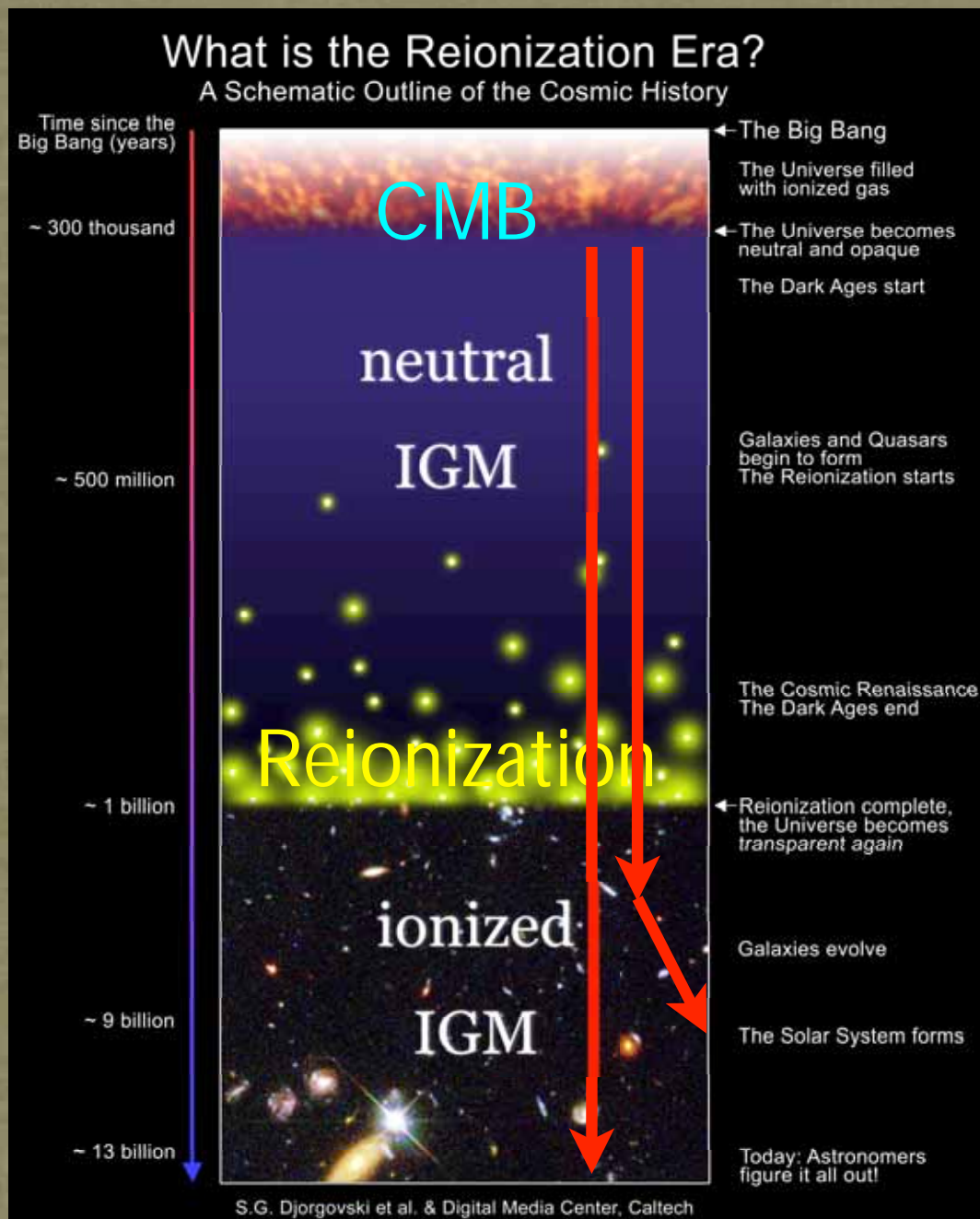
Scales smaller than diffusion scale are damped out

diffusion length sensitive to baryon density

“Silk damping tail”



REIONIZATION AKA FOG



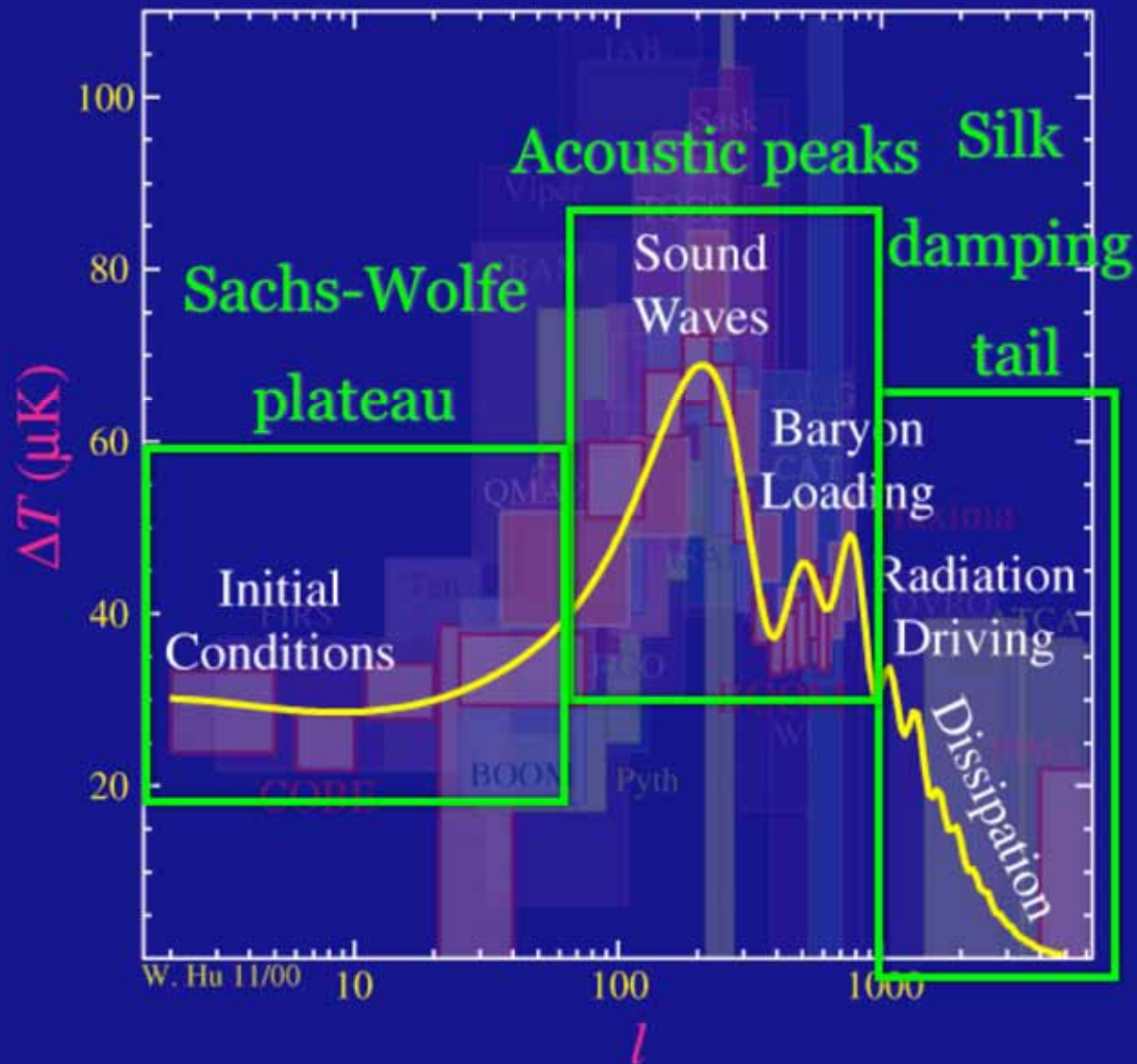
- Galaxies reionize the Universe at $z < 15$
- Ionized IGM rescatters CMB photons
=> damps CMB anisotropies

$$C_l \rightarrow C_l e^{-2\tau}$$

- Optical depth to surface of last scattering $\tau \sim 0.09$

$$\tau = \int \sigma_T n_e(z) dl$$

PRIMARY CMB PHYSICS



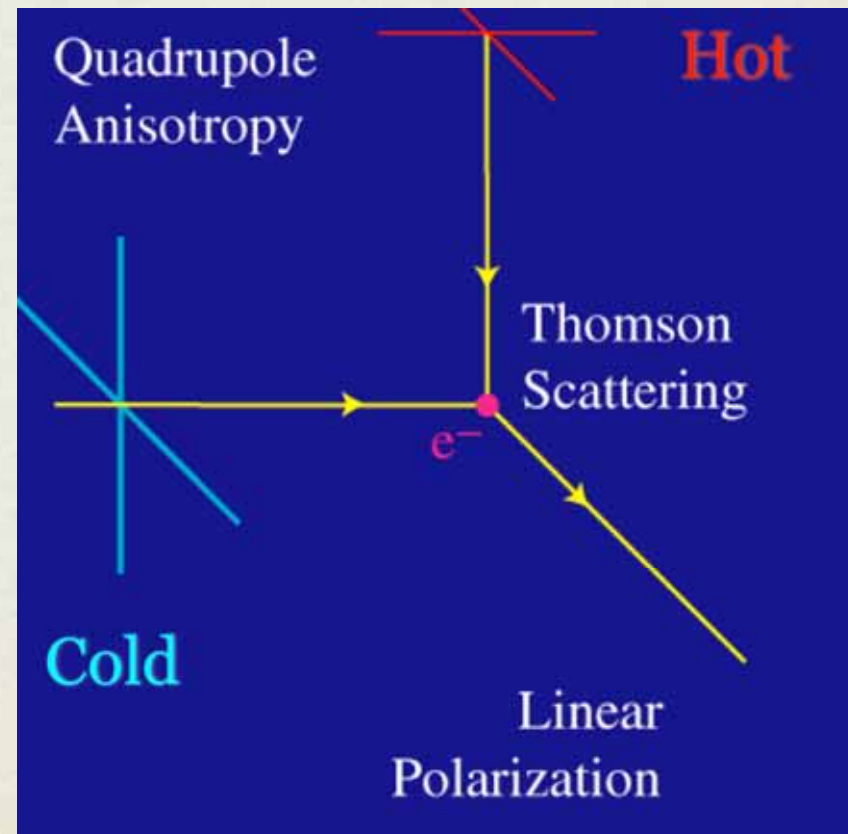
6 numbers to get CMB temperature C_l

Parameter	Planck	
	Best fit	68% limits
$\Omega_b h^2$	0.022068	0.02207 ± 0.00033
$\Omega_c h^2$	0.12029	0.1196 ± 0.0031
$100\theta_{\text{MC}}$	1.04122	1.04132 ± 0.00068
τ	0.0925	0.097 ± 0.038
n_s	0.9624	0.9616 ± 0.0094
$\ln(10^{10} A_s)$	3.098	3.103 ± 0.072

POLARISATION

- Full description of CMB is intensity plus polarisation
- Polarisation originally dismissed as being too small to detect

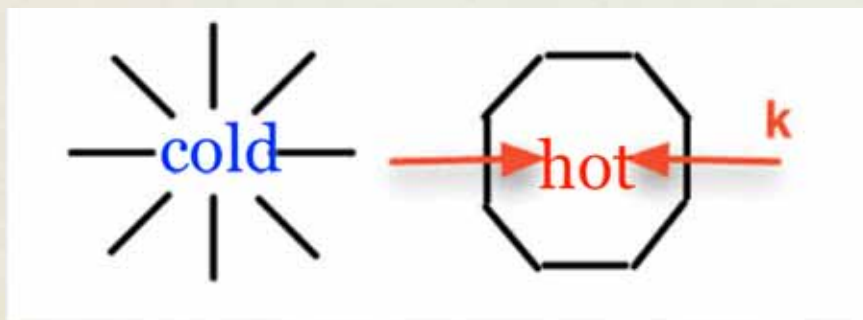
Thompson scattering of
temperature quadrupole
=> linear polarisation



E- AND B-MODE POLARISATION

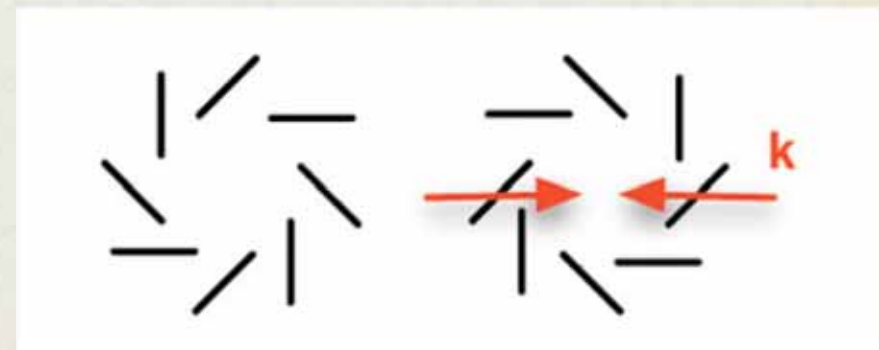
- Density (scalar) fluctuations lead to E-mode polarisation only
- Gravitational waves (tensor) fluctuations lead to B-mode polarisation
- Also get B-mode polarisation by lensing E-mode polarisation

E-mode (grad)

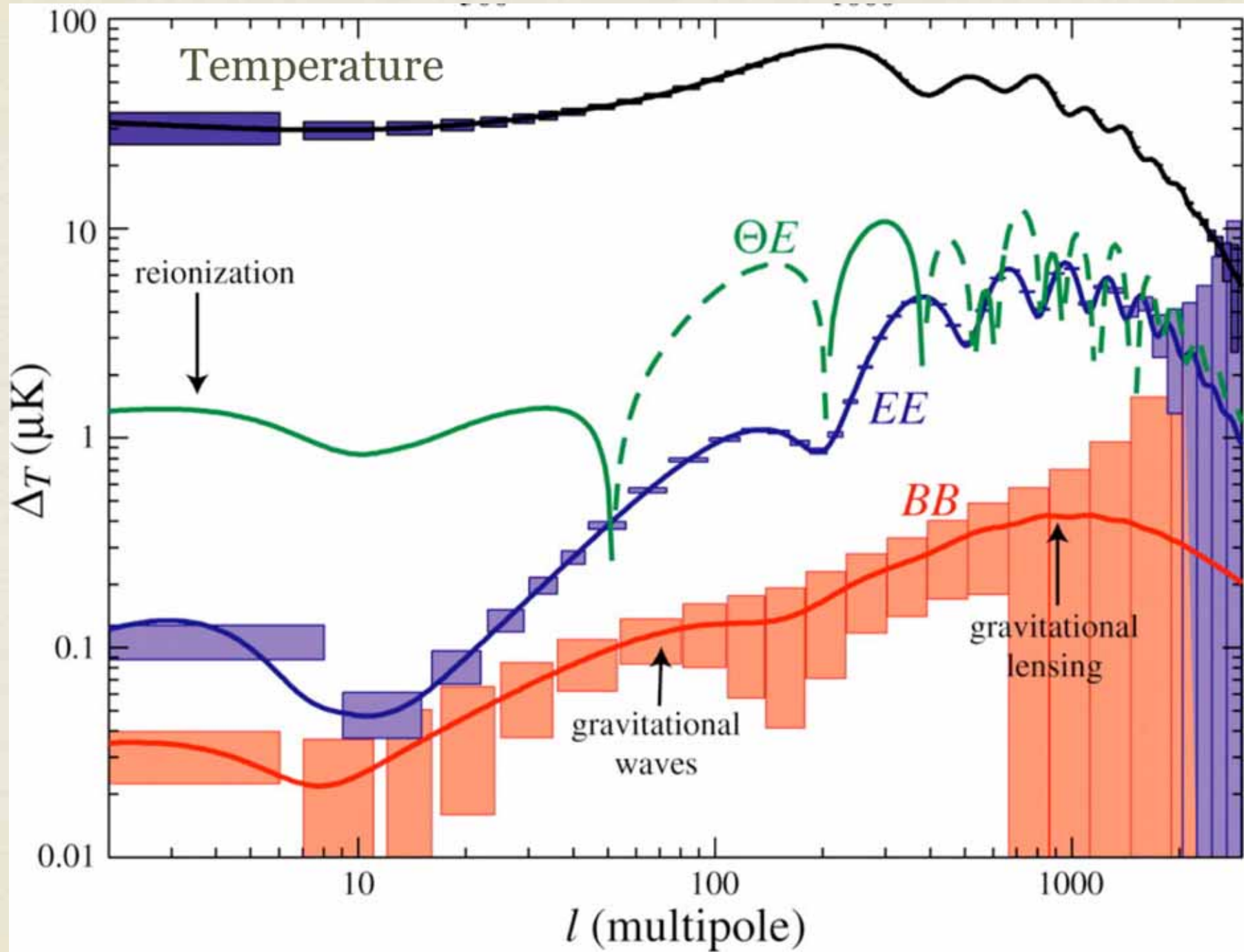


(parity symmetric
=no handedness)

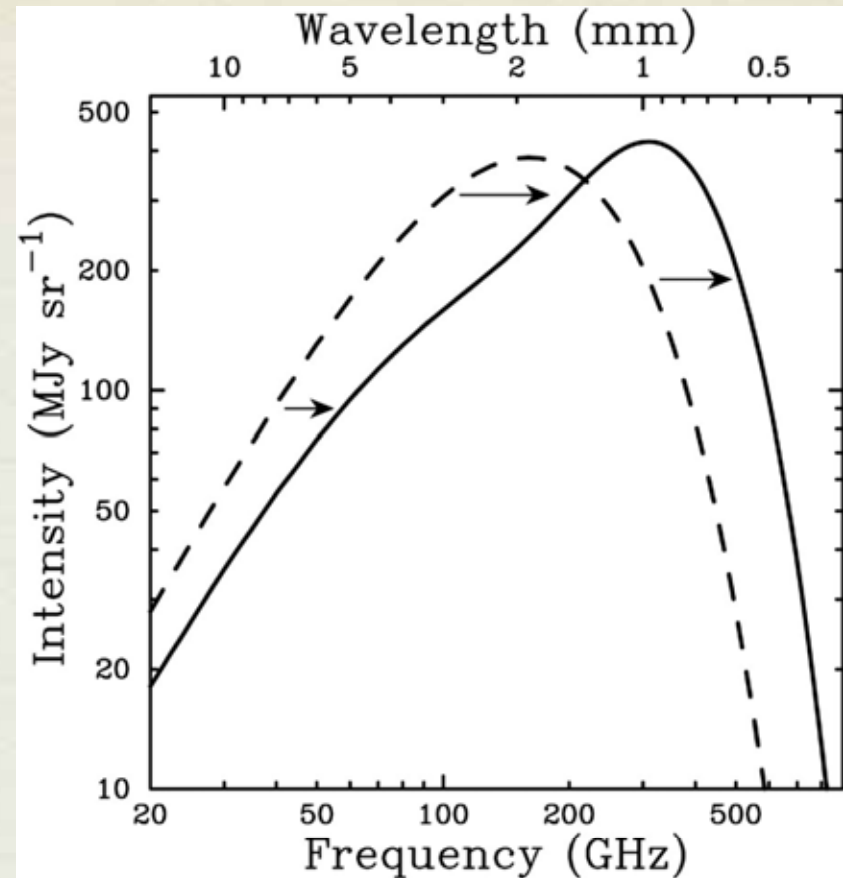
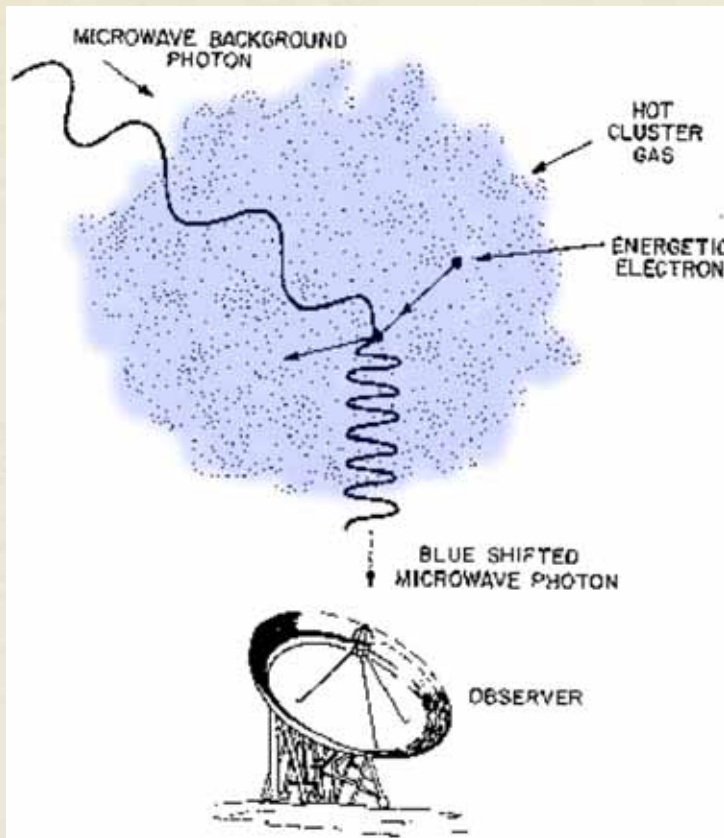
B-mode (curl)



(parity inverted
=handedness)

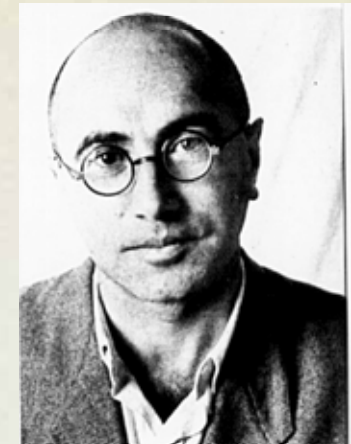
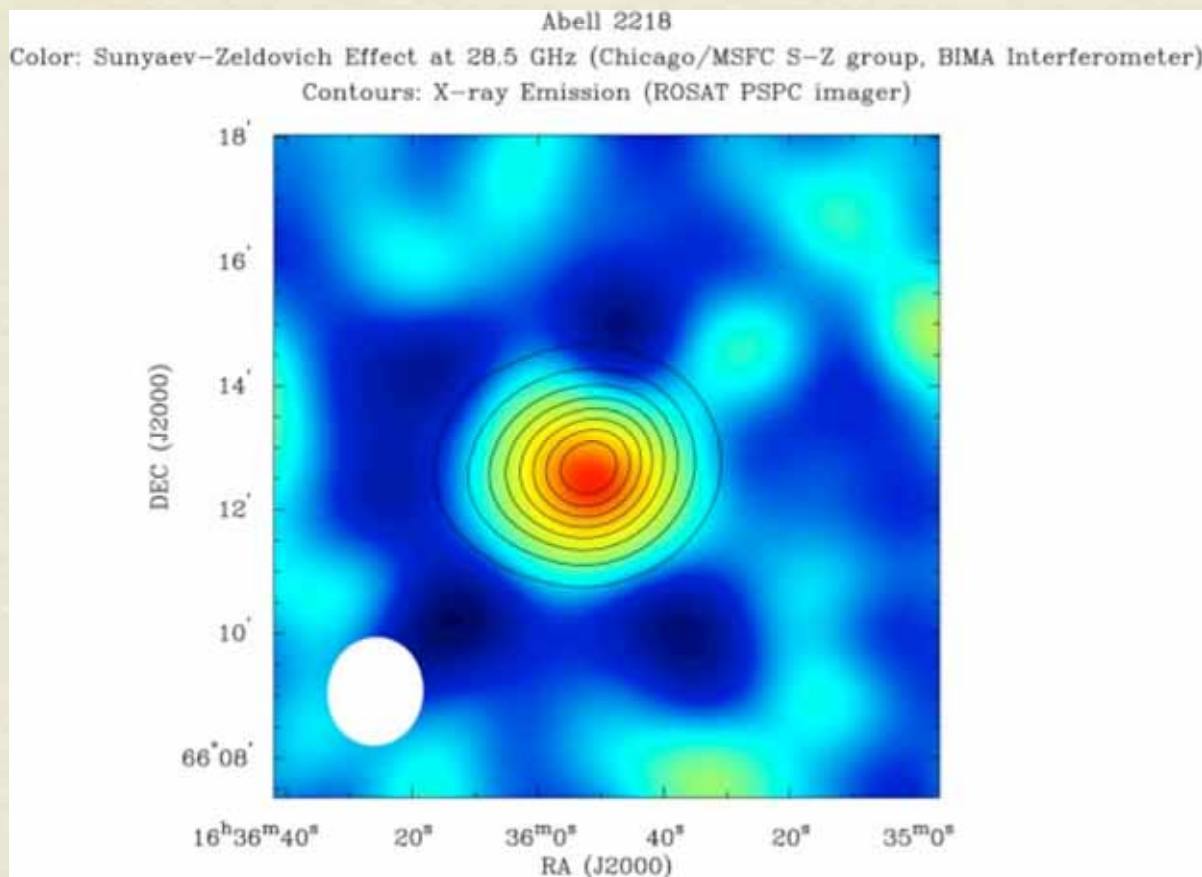


SUNYAEV-ZELDOVICH EFFECT



$$\frac{\Delta T}{T} \approx -2 \int d\tau \frac{k_B T}{m_e c^2}$$

CLUSTERS



Clusters are filled with hot (10^7 - 10^8 K) gas

=> SZ effect makes them show up as small hot spots (~ 10 arcsec)

TSZ VS KSZ

Thermal SZ

$$\frac{\Delta T}{T} \approx -2 \int d\tau \frac{k_B T}{m_e c^2}$$

Thermal velocities

Frequency dependent

Kinetic SZ

$$\frac{\Delta T}{T} = \int d\tau \frac{\mathbf{v} \cdot \hat{\mathbf{n}}}{c}$$

Bulk velocities

Frequency independent

KSZ typically smaller

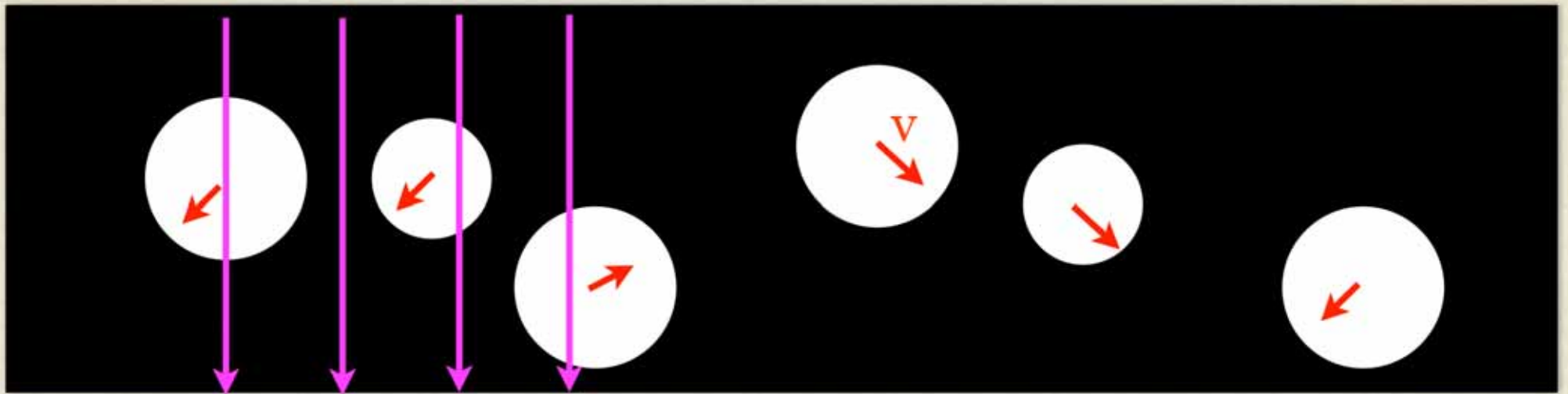
$$\frac{\Delta T_{\text{KSZ}}}{\Delta T_{\text{TSZ}}} \sim \frac{v_r/c}{\langle v_e^2 \rangle / c^2} \sim 1\%$$

cluster peculiar velocity < 1000 km/s

electron thermal velocity $\sim 10^4$ km/s for 10^7 K

Rapaeli 2005

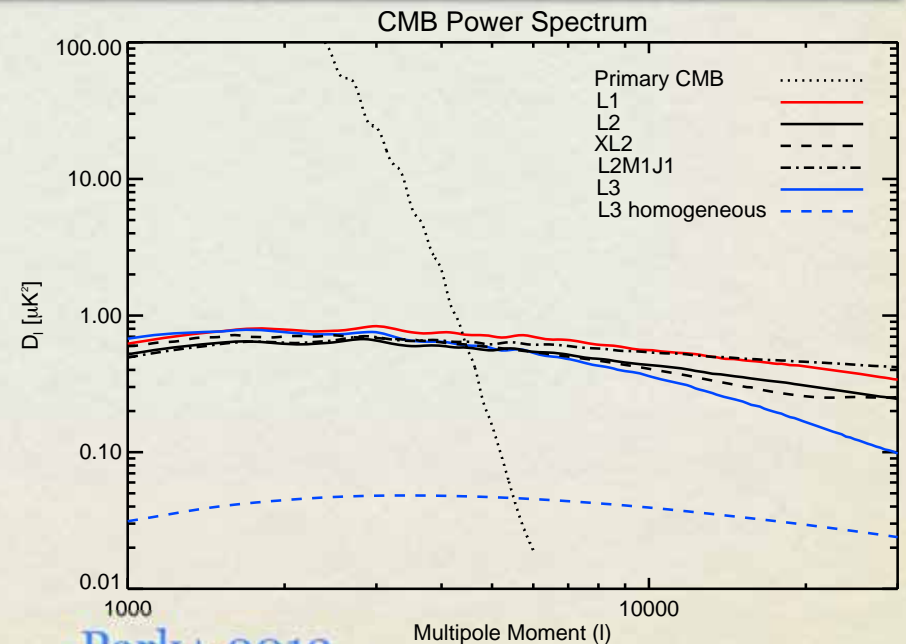
PATCHY KSZ



Bulk velocity of ionized bubbles
introduces patchy KSZ effect
=> probe of reionization

$$\frac{\Delta T_{\text{ksZ}}}{T_{\text{CMB}}}(\hat{\mathbf{n}}) = \frac{\sigma_T}{c} \int_{z_1}^{z_2} \frac{dx}{dz} \frac{dz}{(1+z)} \bar{n}_e(z) e^{-\tau(z)} \hat{\mathbf{n}} \cdot \mathbf{q},$$

$$\mathbf{q} = (1 + \delta_x)(1 + \delta_b) \mathbf{v}.$$

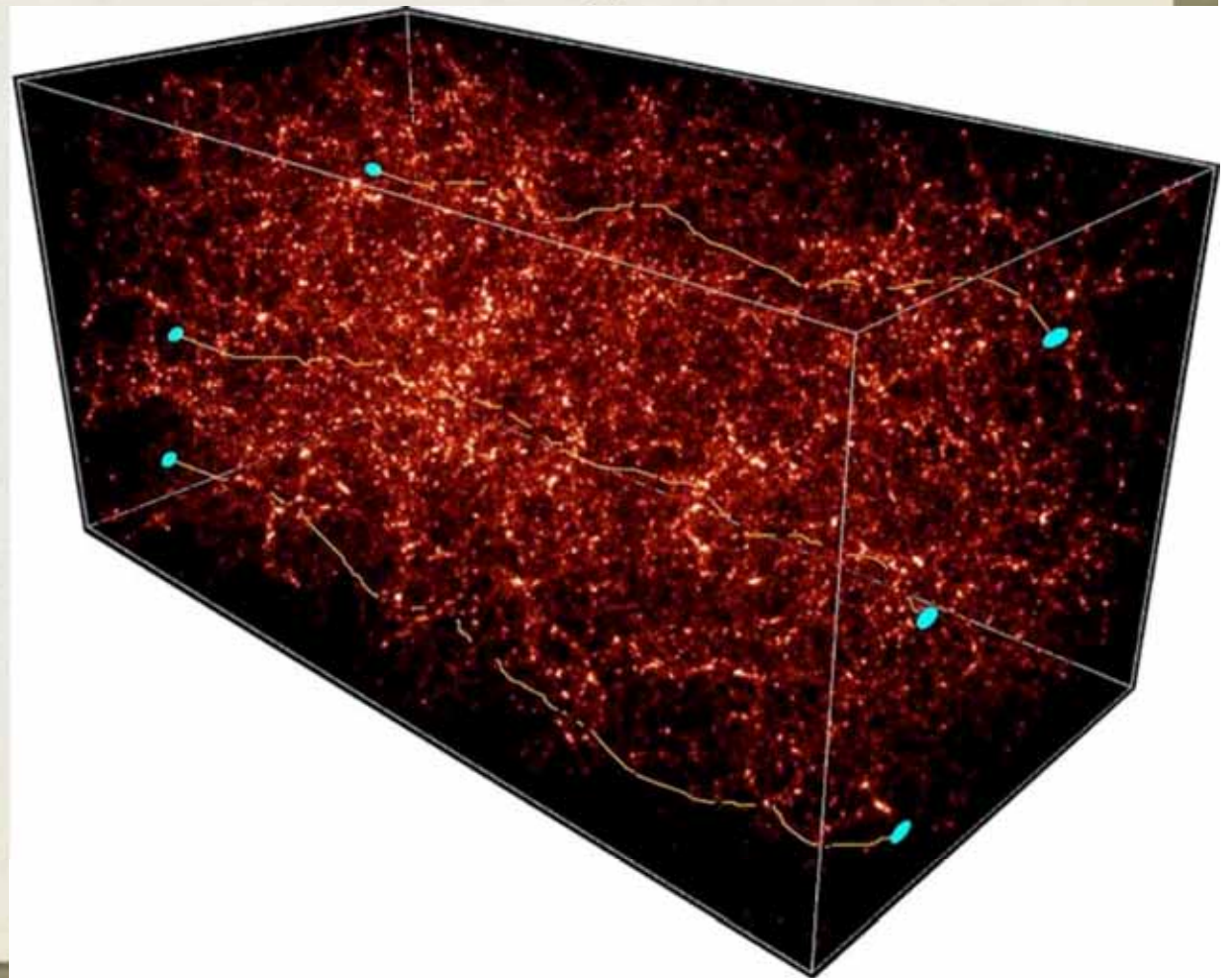


WEAK LENSING

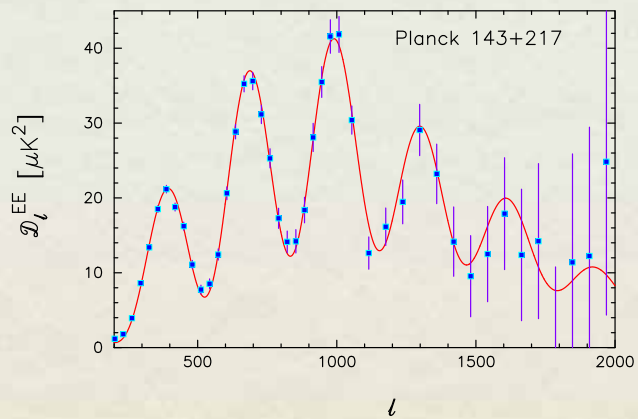
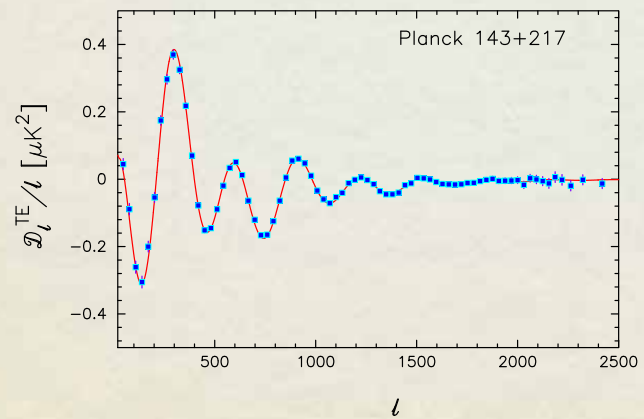
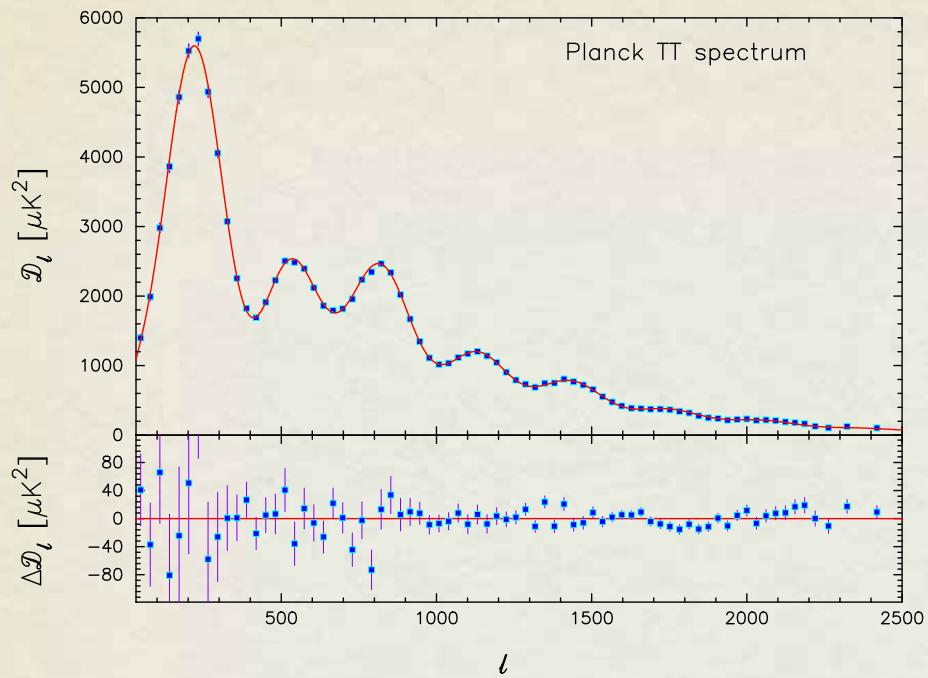
$$T_{\text{lensed}}(\hat{n}') \approx T_{\text{unlensed}}(\hat{n}) + \alpha \cdot \nabla T_{\text{unlensed}}(\hat{n})$$

deflection angle $\alpha = \nabla\phi$

- Matter along line of sight to CMB will gravitationally lens temperature fluctuations
- 1) Mixes E \leftrightarrow B pol
2) Adds Non-Gaussianity
- CMB a background screen to statistically probe large scale structure



PLANCK SPECTRA



PLANCK SCIENCE

- Independent check of WMAP observations
- Inflation models: scalar tilt + better B-mode constraints
- Bispectrum - Non-Gaussianity
- Weak lensing - map matter
- SZ effect - Clusters